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Ball Box System in \mathbb{Z}

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There is a box at each integer $x \in \mathbb{Z}$.

Ball configuration $\eta \in \{0, 1\}^{\mathbb{Z}}$: 0 = empty box, 1 = ball.

Finite number of balls:

Empty carrier starts to the left of leftmost ball, and visits boxes from left to right.

Let $T\eta$: configuration after the carrier visited all boxes.

0	0	1	0	1	1	0	0	0	1	1	1	0	1	0	0	0	0	0	η	
	0	0	1	0	1	2	1	0	0	1	2	3	2	3	2	1	0	0	0	carrier
0	0	0	1	0	0	1	1	0	0	0	0	1	0	1	1	1	0	0	$T\eta$	

Dynamical system, nothing random in the evolution.

Ball-Box-System introduced by [Takahashi-Satsuma \(1990\)](#)

KdV equation

BBS has solitons, a phenomenon present in the Korteweg & de Vries (KdV) differential equation for $u(r, t) \in \mathbb{R}^+$, $r \in \mathbb{R}$, $t \in \mathbb{R}^+$ given by

$$\dot{u} = u''' + u u' \quad (1)$$

For the relation between BBS and KdV see Tokihiro et al , Takahashi and Matsukidaira and Kato, Satoshi and Zuk .

But the way is not easy: From the paper of Tokihiro et al:

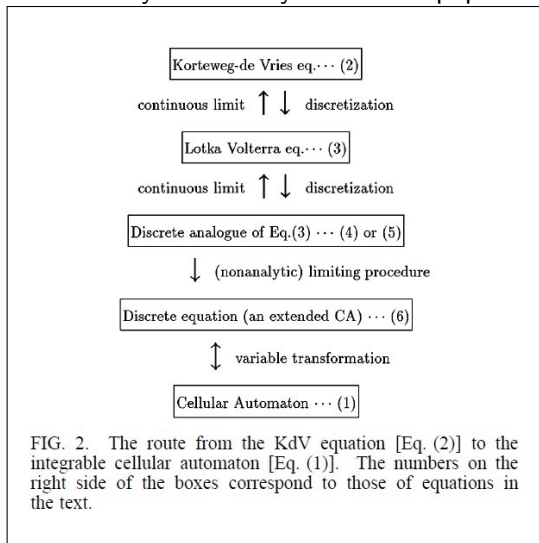


FIG. 2. The route from the KdV equation [Eq. (2)] to the integrable cellular automaton [Eq. (1)]. The numbers on the right side of the boxes correspond to those of equations in the text.

Here BBS with **infinitely many balls**.

If there are “more balls than empty boxes”, the carrier gets an increasing number of balls and the result is just a global flipping:

$T\eta(x) = 1 - \eta(x)$. Hence consider configurations with “upper density less than half”:

$$\mathcal{X}_{\frac{1}{2}} := \left\{ \eta \in \{0, 1\}^{\mathbb{Z}} : \limsup_{n \rightarrow \pm\infty} \frac{1}{n} \left| \sum_{x=0}^n \eta(x) \right| < 1/2 \right\},$$

If $\eta \in \mathcal{X}_{\frac{1}{2}}$, then the carrier load after visiting box x defined by

$$C(x, \eta) := \left(\sup_{z \leq x} \sum_{y=z}^x (2\eta(y) - 1) \right)^+$$

is bounded and $(T\eta)(x) := (C(x-1, \eta) - C(x, \eta))^+$ is well defined.

Since $T\eta$ may not belong to $\mathcal{X}_{\frac{1}{2}}$ we define

$$\mathcal{X} := \left\{ \eta \in \mathcal{X}_{\frac{1}{2}} : \limsup_{y \rightarrow \pm\infty} \left| \frac{1}{y} \sum_{x=0}^y C(x, \eta) \right| < \infty \right\}.$$

We show that if $\eta \in \mathcal{X}$, then $T^n\eta \in \mathcal{X}_{\frac{1}{2}}$ and the dynamics is well defined for all times.

The operator $T : \mathcal{X} \rightarrow \mathcal{X}$ induces operators in the space of bounded functions: $(Tf)(\eta) = f(T\eta)$ and of measures: $(\mu T)f = \mu(Tf)$.

Invariant measures A measure is invariant if $\mu T = \mu$.

Let ν_λ product with density λ (iid Bernoulli(λ)).

Theorem 1 ν_λ is invariant for T for all $\lambda < \frac{1}{2}$.

Proof Application of Burke Theorem. Under iid Bernoulli(λ), carrier performs a **reflecting random walk** on $\{0, 1, 2, \dots\}$ with

$$p(\ell, \ell + 1) = \lambda, p(\ell + 1, \ell) = 1 - \lambda \text{ for } \ell \geq 0 \text{ and } p(0, 0) = 1 - \lambda.$$

Since $\lambda < 1/2$, Geometric($\lambda/(1 - \lambda)$) is **reversible**.

By reversibility, up-jumps \sim reflected down-jumps.

up-jumps = balls in η , down-jumps = balls in $T\eta$.

iid Bernoulli reflecting invariant implies down jumps \sim iid Bernoulli(λ).

□

Goal Characterize the set of shift-ergodic invariant measures for T .

Conserved quantities Number of balls is conserved for $\eta \in \mathcal{X}$.

Number of $\eta(x)(1 - \eta(x + 1))$ is conserved.

We call *k-pode* a set of k successive ones followed by k zeroes in the middle of zeroes. k -podes are conserved and travel at speed k :

```
00000111000000000000000000000000
00000000011100000000000000000000
00000000000001110000000000000000
000000000000000001110000
```

Isolated k -podes travel at speed k and conserve the distances:

```
0000011100000000000000000000000001110000000000000000000000000000000000000000
0000000001110000000000000000000000000000000000000000111000000000000000000000
00000000000001110000000000000000000000000000000000000000000000000000000000000000
00000000000000000111000000000000000000000000000000000000000000000000000000000000
```


k -podes are conserved even when interacting with m -podes:

0000011100000001000000000001110000000000000000000000
00000000011100000100000000000000111000000000000000000000
00000000000011100010000000000000001110000000000000000000
0000000000000001110100000000000000000000001110000000000000
00000000000000000101110000000000000000000011100000000000
000000000000000000100011100000000000000000011100000000

Method of TS to identify k -podes (basic sequences) when they are in the same cycle.

TS algorithm to identify k -podes when there is a finite number of balls

Run: sequence of successive zeroes between two ones or sequence of ones between two zeroes.

Start from the left and look at the runs from left to right.

When a run has length $k < m$, the length of the following run, then the k elements of the short run and the first k elements of the long run are called k -pode.

Identified k -podes are ignored by the algorithm.

Repeat the procedure until all multipodes have been identified.

Records Given $\eta \in \mathcal{X}$, we say that there is a *record* at x if

$$\sum_{y=x-z}^x (2\eta(y) - 1) < 0, \quad \text{for all } z \geq 0.$$

Trajectory of a nearest neighbor walk $(\xi(x) : x \in \mathbb{Z})$ with increments

$$\xi(x) - \xi(x - 1) = 2\eta(x) - 1$$

Record at x for the walk if $\xi(x) < \xi(x - z)$ for all $z > 0$.

Use dots for records:

```

.....111000.....10.....111000.....
.....111000..10.....111000.....
.....11100010.....111000.....
.....11101000.....111000.....
.....10111000.....111000.....
.....10..111000.....111000.....

```

```
00...1100...1100...1100...11001100...1100...1010...10...10...10...10...
1100...1100...1100...1100...11001100...1100...1010...10...10...10...10...
...1100...1100...1100...1100...11001100...1100...1010...10...10...10...10...
...1100...1100...1100...1100...11001100...11001010...10...10...10...10...
...1100...1100...1100...1100...11001100...11010100...10...10...10...10...
...1100...1100...1100...1100...11001100...1010110010...10...10...10...
...1100...1100...1100...1100...11001100...1010110100...1010...10...10...
...1100...1100...1100...1100...11001100101010...10...110010...10...10...
...1100...1100...1100...1100...110011010100...10...110100...10...10...
...1100...1100...1100...1100...11001010110010...10110100...10...
...1100...1100...1100...1100...11001100101000...10...101100...10...
...1100...1100...1100...1100...10101100101100...10...10...1100...10...
...1100...1100...1100...1100...110010101101001010...110010...
...1100...1100...1100...1101010010110010110100...110010...
...1100...1100...1100...10101010...110100...110100101100...101100...
...1100...1100...11001010...101100...101101001100...10...1100...
...1100...1100...11010100...10...110010...1011001100...10...1100...
...1100...1100...1010110010...11010010...1100110010...1100...
...1100...1100...1010...110100...10110100...11001100...1100...
...1100...1100...1010...101100...10...1100...10...1100...11001100...1100...
...1100...1100...1010...101100...10...1100...10...1100...11001100...1100...
...11001010...10...110010...10...110010...11001100100...1100...
...1101000...10...110010...110100...110011001100...1100...
...1010110010...10110100...101100...1100110011001100...1100...
...1010110100...10...101100...10...1100...110011001100...1100...
...1010...101100...10...10...1100...10...1100...110011001100...1100...
...1010...10...110010...10...1100...10...1100...110011001100...1100...
...1010...10...11010010...110010...1100...110011001100...1100...
...1010...10...10110100...110100...1100...110011001100...1100...
...1010...10...10...101100...101100...1100...110011001100...1100...
...1010...10...10...10...1100...10...1100...110011001100...1100...
...1010...10...10...10...110010...1100...1100...110011001100...1100...
...1010...10...10...10...110010...1100...1100...110011001100...1100...
...1010...10...10...10...110100...1100...1100...110011001100...1100...
...1010...10...10...10...101100...1100...1100...110011001100...1100
```

Effective distance between successive k -podes

Number of records between them when isolated from m -podes for $m > k$.

```
.....111000.....10.....10.....11011000.....
.....111000.....10.....10.....10111000.....
.....11100010.....10.....10.....111000.....
.....1110100010.....10.....10.....111000.....
.....1011100100.....10.....111000.....
.....10.....11011000.....10.....111000.....
.....10.....1011100010.....111000.....
```

For $m > k$:

- When m -pode overpasses k -pode, the k -pode gains $2(m - k)$ records with respect to its natural trajectory.
- k -podes do not affect effective distance of m -podes.

The tagged record process System as seen from a tagged record.

Configurations with a record at the origin:

$$\mathcal{X}^o := \eta \in \mathcal{X} : \text{there is a record at } 0\}$$

Take $\eta \in \mathcal{X}^o$,

$$r_t(\eta, i) := \text{position in } T^t\eta \text{ of } i\text{th record of } \eta$$

```
.....111000..10...10..x...111000.....
.....11100010...10.x...111000.....
.....11101000.10x.....111000.....
.....x.1011100100.....111000.....
.....x.10..11011000.....111000.....
.....x..10...10.111000.....111000..
```

x is the tagged record $r_t(\eta, 0)$

For $\eta \in \mathcal{X}^o$ system at time t as seen from the tagged record:

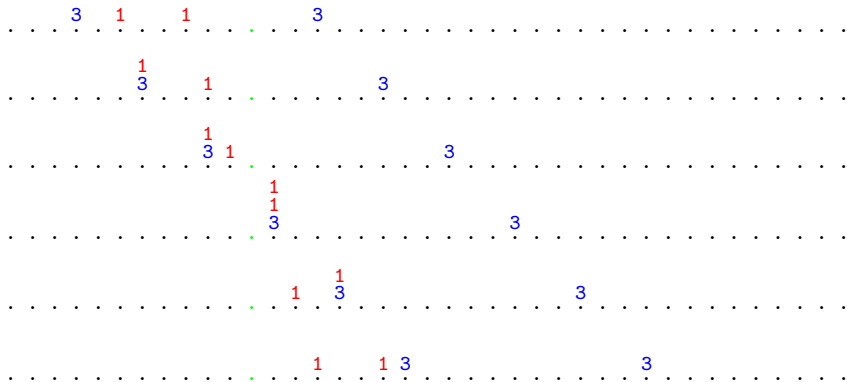
$$\hat{T}^t \eta = \theta_{r_t(\eta, 0)} T^t \eta$$

```

..... 111000 . 10 . . 10 . x . . 111000 . . . . .
..... 11100010 . . 10 . x . . . 111000 . . . . .
..... 11101000 . 10x . . . . . 111000 . . . . .
..... . . . . . . . . . . x . 1011100100 . . . . . 111000 . . . . .
..... . . . . . . . . . . x . 10 . 11011000 . . . . . 111000 . . . . .
..... . . . . . . . . . . x . 10 . . 10 . 111000 . . . . . 111000 . . . . .

```

Fuzzy representation



3-pode motion is predictable from the single time fuzzy configuration.

The 1-pode evolution can be reconstructed from the last 3 steps.

Shift-ergodicity Shift $(\theta_x \eta)(y) = \eta(y - x)$.

A shift-stationary μ is *ergodic* if for all measurable $f : \mathcal{X} \rightarrow \mathbb{R}$,

$$\theta_x f = f \quad \text{implies } f \text{ is constant } \mu\text{-a.s.}$$

Equivalent to Cesaro asymptotic independence:

$$\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{x=0}^{n+1} f(\eta)g(\theta_x \eta) = \mu f \mu g.$$

for continuous f and g .

Record-shift-ergodicity μ is *record-stationary* if $\mu(\mathcal{X}^o) = 1$ and

$$\int f(\theta_{r(\eta,i)}\eta)\mu(d\eta) = \mu f$$

for all record i and measurable f . (point stationarity, Thorisson).

A record-stationary μ is *record-ergodic* if for all $f : \mathcal{X}^o \rightarrow \mathbb{R}$

$$\theta_{r(\eta,i)}f(\eta) = f(\eta), \text{ for all } \eta, i \quad \text{implies } f \text{ is constant } \mu\text{-a.s.}$$

Equivalent to Cesaro asymptotic independence along records:

$$\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=0}^{n+1} f(\eta)g(\theta_{r(\eta,i)}\eta) = \mu f \mu g, \quad \mu\text{a.s.}$$

for continuous f and g .

Palm measure μ shift-stationary.

$\hat{\mu} := \mu$ conditioned to record at origin:

$$\hat{\mu}f = \frac{1}{\mu(\mathcal{X}^o)} \int f(\eta) \mathbf{1}\{\eta \in \mathcal{X}^o\} \mu(d\eta).$$

Proposition 2 (a) μ shift-stationary iff $\hat{\mu}$ record-shift stationary.

(b) μ ergodic iff $\hat{\mu}$ record-ergodic.

Theorem 3 (Harris, Port & Stone) Let μ be shift-stationary. Then,

$$\hat{\mu}\hat{T} = (\mu T)^\wedge$$

In particular $\mu T = \mu$ if and only if $\hat{\mu}\hat{T} = \hat{\mu}$.

Proof uses that T commutes with shift: $T\theta_x = \theta_x T$, for all $x \in \mathbb{Z}$.

Multipode decomposition

Family of operators $M_k : \mathcal{X}^o \rightarrow \mathcal{X}^o$.

$M_k \eta$ has only k -podes at its *effective distance*.

.....11101000.10x..10...111000.10.....	η
.....10...10x..10.....10.....	$M_1 \eta$
.....x.....	$M_2 \eta$
.....111000.x.....111000.....	$M_3 \eta$
...	...

Effective distance between successive k -podes computed by evolving the piece of configuration containing the cycles of both k -podes and everything between them.

Define $M : \mathcal{X}^o \rightarrow (\mathcal{X}^o)^{\mathbb{N}}$ given by $M : \eta \mapsto (M_k \eta : k \geq 1)$.

Proposition 4 M is one-to-one.

Shift-ergodic invariant measures

Our main result:

Theorem 5 *Let μ be shift-ergodic invariant with $\lambda < 1/4$ and let $\hat{\mu}$ be its Palm measure. Then,*

$$\hat{\mu}M = \bigotimes_{k \geq 1} \hat{\mu}M_k.$$

That is, if η has law $\hat{\mu}$,

$(M_k\eta : k \geq 1)$ is a family of independent configurations.

Recall that $M_k\eta$ contains only k -podes.

Evolution of decomposed configurations

Multipode decomposition of the configuration at time t consists of distinct translations of the decomposition of the initial configuration

Let $y_t(\eta, k, i)$ be the increment in records of the first k -pode of $\hat{T}^t\eta$ to the right of $r(\hat{T}^t\eta, i)$ in the time interval $[0, t]$.

Theorem 6 For $\eta \in \mathcal{X}^o$ and $k \geq 1$ we have

$$M_k \hat{T}^t \eta = M_k \theta_{-y_t(\eta, k)} \eta$$

Let ρ_k be the mean number of k -podes per record under $\hat{\mu}$. Since a k -pode occupies $2k$ sites in a cycle, we have that

$a := 2 \sum_k k \rho_k$ is the mean cycle size under $\hat{\mu}$.

By ergodicity, the density of balls $\lambda = \frac{a/2}{a+1}$. Hence,

$$a = \sum_k 2k \rho_k = \frac{2\lambda}{1 - 2\lambda} < \infty.$$

Proposition 7 μ shift-ergodic with $\lambda < 1/4$ and $\hat{\mu}$ its Palm measure. Assume $\rho_k > 0$, $\rho_\ell > 0$ and $\ell < k$. Then

$$0 < c \leq \underline{\lim}_{t \rightarrow \infty} \frac{y_t(\eta, k, i) - y_t(\eta, \ell, i)}{t} \leq \overline{\lim}_{t \rightarrow \infty} \frac{y_t(\eta, k, i) - y_t(\eta, \ell, i)}{t} \leq k - \ell, \quad (2)$$

for $\hat{\mu}$ almost all η ; $c = c(k, \ell)$.

For the moment the inequality $0 < c$ holds under $\lambda < 1/4$.

Proof of Theorem 5

$$\begin{aligned}\hat{\mu}(M_k f M_m g) &= \int \hat{\mu}(d\eta) f(M_k \hat{T}^t \eta) g(M_m \hat{T}^t \eta) \quad (\text{invariance of } \hat{\mu}) \\ &= \int \hat{\mu}(d\eta) f(M_k \theta_{-y_t(\eta, k, 0)} \eta) g(M_m \theta_{-y_t(\eta, m, 0)} \eta) \quad (\text{Thm 6}) \\ &= \int \hat{\mu}(d\eta) f(M_k \eta) g(M_m \theta_{-y_t(\eta, m, 0) + y_t(\eta, k, 0)} \eta) \quad (\text{stationarity})\end{aligned}$$

Hence,

$$\begin{aligned}\hat{\mu}(M_k f M_m g) &= \int \hat{\mu}(d\eta) \frac{1}{t} \sum_{s=0}^t f(M_k \eta) g(M_m \theta_{-y_s(\eta, m, 0) + y_s(\eta, k, 0)} \eta) \\ &\xrightarrow[t \rightarrow \infty]{} \int \hat{\mu}(d\eta) f(M_k \eta) \int \hat{\mu}(d\eta) g(M_m \eta) \quad (\text{ergodicity of } \hat{\mu}) \\ &= \hat{\mu} M_k f \hat{\mu} M_m g\end{aligned}$$

The limit uses that for large t ,

$$y_t(\eta, m, 0) - y_t(\eta, k, 0) \geq tc(k, m) \rightarrow \infty. \quad (\text{Proposition 7, b}) \quad \square$$

Speed of k -podes

Proposition 8 μ shift-ergodic and $\hat{\mu}$ its Palm measure. Assume $\rho_k > 0$ and $\sum_{m \geq 1} m^2 \rho_m < \infty$. Then the following limits exist

$$\lim_{t \rightarrow \infty} \frac{y_t(\eta, k, i)}{t} = v_k < \infty \quad (3)$$

for $\hat{\mu}$ almost all η . The limits are the unique solution of the system

$$v_k = k + \sum_{m=k+1}^{\infty} 2(m-k)(v_m - v_k)\rho_m, \quad k \geq 1. \quad (4)$$

and satisfy the following bounds:

$$\frac{\ell}{c_\ell} + 2 \sum_{m>\ell} m(m-\ell) \frac{\rho_m}{c_m} \leq v_\ell \leq \frac{\ell}{c_\ell} + 2c_\ell \sum_{m>\ell} m(m-\ell) \frac{\rho_m}{c_m} \quad (5)$$

References

- [1] T. E. Harris. Random measures and motions of point processes. *Z. Wahrscheinlichkeitstheorie und Verw. Gebiete*, 18:85–115, 1971.
- [2] T. Kato, S. Tsujimoto, and A. Zuk. Spectral analysis of transition operators, automata groups and translation in BBS. *Communications in Mathematical Physics*, pages 1–25, 2016.
- [3] T. M. Liggett. *Interacting particle systems*. Classics in Mathematics. Springer-Verlag, Berlin, 2005. Reprint of the 1985 original.
- [4] J. W. Pitman. One-dimensional Brownian motion and the three-dimensional Bessel process. *Advances in Appl. Probability*, 7(3):511–526, 1975.
- [5] S. C. Port and C. J. Stone. Infinite particle systems. *Trans. Amer. Math. Soc.*, 178:307–340, 1973.
- [6] D. Revuz and M. Yor. *Continuous martingales and Brownian motion*, volume 293 of *Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences]*. Springer-Verlag,

Berlin, third edition, 1999.

- [7] D. Takahashi and J. Matsukidaira. Box and ball system with a carrier and ultradiscrete modified kdv equation. *Journal of Physics A: Mathematical and General*, 30(21):L733, 1997.
- [8] D. Takahashi and J. Satsuma. A soliton cellular automaton. *Journal of The Physical Society of Japan*, 59(10):3514–3519, 1990.
- [9] H. Thorisson. *Coupling, stationarity, and regeneration*. Probability and its Applications (New York). Springer-Verlag, New York, 2000.
- [10] H. Thorisson. *Coupling, stationarity, and regeneration*, volume 200. Springer New York, 2000.
- [11] T. Tokihiro, D. Takahashi, J. Matsukidaira, and J. Satsuma. From soliton equations to integrable cellular automata through a limiting procedure. *Physical Review Letters*, 76(18):3247, 1996.