"Mostly Maximum Principle" 5th edition: in Latin America for the first time

# A Schiffer-type problem in annuli and applications to Euler flows

# Pieralberto Sicbaldi

Joint work with A. Enciso, A. J. Fernández and D. Ruiz

Rio de Janeiro

June 2024

To find domains  $\Omega \subset \mathbb{R}^2$  that support a solution of the overdetermined problem

$$\begin{cases} \Delta u + \mu u &= 0 & \text{in } \Omega \\ \frac{\partial u}{\partial \nu} &= 0 & \text{on } \partial \Omega \\ u &= \text{ constant } \text{on } \partial \Omega \end{cases}$$

where  $\mu$  is a constant and  $\nu$  is the normal vector on  $\partial\Omega$ .

To find domains  $\Omega \subset \mathbb{R}^2$  that support a solution of the overdetermined problem

$$\begin{cases} \Delta u + \mu u &= 0 & \text{in } \Omega \\ \frac{\partial u}{\partial \nu} &= 0 & \text{on } \partial \Omega \\ u &= \text{ constant } \text{on } \partial \Omega \end{cases}$$

where  $\mu$  is a constant and  $\nu$  is the normal vector on  $\partial\Omega$ .

Schiffer conjecture ('50 - stated by Yau in 1982). If  $\Omega$  is bounded and smooth and u is nonconstant, then  $\Omega$  is a disk.

To find domains  $\Omega \subset \mathbb{R}^2$  that support a solution of the overdetermined problem

$$\begin{cases} \Delta u + \mu u &= 0 & \text{in } \Omega \\ \frac{\partial u}{\partial \nu} &= 0 & \text{on } \partial \Omega \\ u &= \text{ constant } \text{on } \partial \Omega \end{cases}$$

where  $\mu$  is a constant and  $\nu$  is the normal vector on  $\partial\Omega$ .

Schiffer conjecture ('50 - stated by Yau in 1982). If  $\Omega$  is bounded and smooth and u is nonconstant, then  $\Omega$  is a disk.

The conjecture is still open.

Let  $f : \mathbb{R}^2 \to \mathbb{R}$  an (unknown) continuous function and  $\Omega \subset \mathbb{R}^2$  a given bounded domain. Is it possible to determine f if we know the values

$$\int_{\mathcal{R}(\Omega)} f(x) \mathrm{d}x$$

for any rigid motion  $\mathcal{R}$  in  $\mathbb{R}^2$ ?

Let  $f : \mathbb{R}^2 \to \mathbb{R}$  an (unknown) continuous function and  $\Omega \subset \mathbb{R}^2$  a given bounded domain. Is it possible to determine f if we know the values

$$\int_{\mathcal{R}(\Omega)} f(x) \mathrm{d}x$$

for any rigid motion  $\mathcal{R}$  in  $\mathbb{R}^2$ ?

Applications: Remote sensing, Image recovery, Tomography.

Let  $f : \mathbb{R}^2 \to \mathbb{R}$  an (unknown) continuous function and  $\Omega \subset \mathbb{R}^2$  a given bounded domain. Is it possible to determine *f* if we know the values

$$\int_{\mathcal{R}(\Omega)} f(x) \mathrm{d}x$$

for any rigid motion  $\mathcal{R}$  in  $\mathbb{R}^2$ ?

Applications: Remote sensing, Image recovery, Tomography.

The answer is YES if and only if  $\Omega$  has the **Pompeiu property**:

$$\left[\int_{\mathcal{R}(\Omega)} f(x) \mathrm{d}x = 0 \ \forall \mathcal{R}\right] \Longrightarrow f \equiv 0$$

Let  $f : \mathbb{R}^2 \to \mathbb{R}$  an (unknown) continuous function and  $\Omega \subset \mathbb{R}^2$  a given bounded domain. Is it possible to determine *f* if we know the values

$$\int_{\mathcal{R}(\Omega)} f(x) \mathrm{d}x$$

for any rigid motion  $\mathcal{R}$  in  $\mathbb{R}^2$ ?

Applications: Remote sensing, Image recovery, Tomography.

The answer is YES if and only if  $\Omega$  has the **Pompeiu property**:

$$\left[\int_{\mathcal{R}(\Omega)} f(x) \mathrm{d}x = 0 \ \forall \mathcal{R}\right] \Longrightarrow f \equiv 0$$

Ellipses and Polygons have the Pompeiu property.

## Equivalence

#### Equivalence

**Williams' Theorem (1976).** If  $\Omega$  is simply connected then  $\Omega$  has the Pompeiu property if and only if the Schiffer problem cannot be solved in  $\Omega$ .

#### Equivalence

**Williams' Theorem (1976).** If  $\Omega$  is simply connected then  $\Omega$  has the Pompeiu property if and only if the Schiffer problem cannot be solved in  $\Omega$ .

S = class of smooth simply connected domains (bounded).

**Williams' Theorem (1976).** If  $\Omega$  is simply connected then  $\Omega$  has the Pompeiu property if and only if the Schiffer problem cannot be solved in  $\Omega$ .

S = class of smooth simply connected domains (bounded).

**Open problem 1.** Is it true that disks are the only domains in  $\mathcal{S}$  not having the Pompeiu property?

**Williams' Theorem (1976).** If  $\Omega$  is simply connected then  $\Omega$  has the Pompeiu property if and only if the Schiffer problem cannot be solved in  $\Omega$ .

S = class of smooth simply connected domains (bounded).

**Open problem 1.** Is it true that disks are the only domains in  $\mathcal{S}$  not having the Pompeiu property?

**Open problem 2.** Is it true that disks are the only domains in S where the Schiffer Problem can be solved? (most important case of the Schiffer conjecture)

1.  $\partial\Omega$  is connected and there exists an infinite sequence of orthogonal Neumann eigenfunctions that are constant on  $\partial\Omega$  (Berenstein 1980, Berenstein-Yang 1987).

1.  $\partial\Omega$  is connected and there exists an infinite sequence of orthogonal Neumann eigenfunctions that are constant on  $\partial\Omega$  (Berenstein 1980, Berenstein-Yang 1987).

2.  $\partial \Omega$  is connected and the third order interior normal derivative of u is constant on  $\partial \Omega$  (Liu, 2010).

1.  $\partial\Omega$  is connected and there exists an infinite sequence of orthogonal Neumann eigenfunctions that are constant on  $\partial\Omega$  (Berenstein 1980, Berenstein-Yang 1987).

2.  $\partial \Omega$  is connected and the third order interior normal derivative of *u* is constant on  $\partial \Omega$  (Liu, 2010).

3.  $\Omega$  is simply connected and u has no saddle points in the interior of  $\Omega$  (Willms-Gladwell, 1994).

1.  $\partial\Omega$  is connected and there exists an infinite sequence of orthogonal Neumann eigenfunctions that are constant on  $\partial\Omega$  (Berenstein 1980, Berenstein-Yang 1987).

2.  $\partial \Omega$  is connected and the third order interior normal derivative of u is constant on  $\partial \Omega$  (Liu, 2010).

3.  $\Omega$  is simply connected and u has no saddle points in the interior of  $\Omega$  (Willms-Gladwell, 1994).

4.  $\Omega$  is simply connected and the eigenvalue  $\mu$  is among its seven lowest Neumann eigenvalues (Aviles 1986, Deng 2012).

1.  $\partial\Omega$  is connected and there exists an infinite sequence of orthogonal Neumann eigenfunctions that are constant on  $\partial\Omega$  (Berenstein 1980, Berenstein-Yang 1987).

2.  $\partial \Omega$  is connected and the third order interior normal derivative of *u* is constant on  $\partial \Omega$  (Liu, 2010).

3.  $\Omega$  is simply connected and u has no saddle points in the interior of  $\Omega$  (Willms-Gladwell, 1994).

4.  $\Omega$  is simply connected and the eigenvalue  $\mu$  is among its seven lowest Neumann eigenvalues (Aviles 1986, Deng 2012).

5. The fourth or fifth order interior normal derivative of u is constant on  $\partial \Omega$  (Kawohl-Lucia, 2020).

Let  $\Omega \subset \mathbb{R}^2$  a smooth domain, and let us denote by  $\Gamma_i$  the connected components of  $\partial \Omega$ . Assume that there exists a solution to the "Schiffer-type" problem:

$$\left\{ egin{array}{lll} \Delta u + \mu \, u &= 0 & ext{in} & \Omega \ & & \displaystyle rac{\partial u}{\partial 
u} &= 0 & ext{on} & \partial \Omega \ & & \displaystyle u &= c_i & ext{on} & \partial \Omega \end{array} 
ight.$$

for some some constants  $c_i$  and a positive constant  $\mu$ .

Let  $\Omega \subset \mathbb{R}^2$  a smooth domain, and let us denote by  $\Gamma_i$  the connected components of  $\partial \Omega$ . Assume that there exists a solution to the "Schiffer-type" problem:

$$\left\{ egin{array}{lll} \Delta u + \mu \, u &= 0 & ext{in} & \Omega \ & & \displaystyle rac{\partial u}{\partial 
u} &= 0 & ext{on} & \partial \Omega \ & & \displaystyle u &= c_i & ext{on} & \partial \Omega \end{array} 
ight.$$

for some some constants  $c_i$  and a positive constant  $\mu$ .

A stronger Schiffer conjecture would be that  $\Omega$  must be a disk or an annulus.

Let  $\Omega \subset \mathbb{R}^2$  a smooth domain, and let us denote by  $\Gamma_i$  the connected components of  $\partial \Omega$ . Assume that there exists a solution to the "Schiffer-type" problem:

$$\left\{ egin{array}{lll} \Delta u + \mu \, u &= 0 & ext{in} & \Omega \ & & \displaystyle rac{\partial u}{\partial 
u} &= 0 & ext{on} & \partial \Omega \ & & \displaystyle u &= c_i & ext{on} & \partial \Omega \end{array} 
ight.$$

for some some constants  $c_i$  and a positive constant  $\mu$ .

A stronger Schiffer conjecture would be that  $\Omega$  must be a disk or an annulus.

"Question". Is it true that  $\Omega$  must be a disk or an annulus?

### Shared rigidity results with the original conjecture

Shared rigidity results with the original conjecture

Theorem 1. (Enciso, Fernández, Ruiz, S. 2023)

Theorem 1. (Enciso, Fernández, Ruiz, S. 2023)

1. If there exists an infinite sequence of orthogonal Neumann eigenfunctions that are constant on each component of  $\partial \Omega$ , then this stronger version of the Schiffer conjecture is true.

Theorem 1. (Enciso, Fernández, Ruiz, S. 2023)

1. If there exists an infinite sequence of orthogonal Neumann eigenfunctions that are constant on each component of  $\partial \Omega$ , then this stronger version of the Schiffer conjecture is true.

2. If the eigenvalue  $\mu$  is among its four lowest Neumann eigenvalues, then again this stronger version of the Schiffer conjecture is true.

Theorem 1. (Enciso, Fernández, Ruiz, S. 2023)

1. If there exists an infinite sequence of orthogonal Neumann eigenfunctions that are constant on each component of  $\partial\Omega$ , then this stronger version of the Schiffer conjecture is true.

2. If the eigenvalue  $\mu$  is among its four lowest Neumann eigenvalues, then again this stronger version of the Schiffer conjecture is true.

In this talk we skip the proof of such two results.

Main result (Enciso, Fernández, Ruiz, S. 2023)

In general, the second conjecture is false.

#### Main result (Enciso, Fernández, Ruiz, S. 2023)

#### In general, the second conjecture is false.

**Theorem 2.** For many integer numbers *l* there exist domains  $\Omega_{l,s} \subset \mathbb{R}^2$  given in polar coordinates by

$$\Omega_{l,s} := \{ (r, \theta) : a_l + s \, b_{l,s}(\theta) < r < 1 + s \, B_{l,s}(\theta) \} \,, \ s \in (-\epsilon, \epsilon) \,,$$

where  $a_l \in (0, 1)$  and  $b_{l,s}$ ,  $B_{l,s}$  are analytic functions of the form

$$b_{l,s}(\theta) = \alpha_l \cos l\theta + \mathcal{O}(s), \qquad B_{l,s}(\theta) = \beta_l \cos l\theta + \mathcal{O}(s),$$

with  $\alpha_l$  and  $\beta_l$  nonzero constants, such that the "Schiffer-type" problem can be solved in  $\Omega_{l,s}$ .

#### Main result (Enciso, Fernández, Ruiz, S. 2023)

#### In general, the second conjecture is false.

**Theorem 2.** For many integer numbers *l* there exist domains  $\Omega_{l,s} \subset \mathbb{R}^2$  given in polar coordinates by

$$\Omega_{l,s} := \{ (r, \theta) : a_l + s \, b_{l,s}(\theta) < r < 1 + s \, B_{l,s}(\theta) \} \,, \ s \in (-\epsilon, \epsilon) \,,$$

where  $a_l \in (0, 1)$  and  $b_{l,s}$ ,  $B_{l,s}$  are analytic functions of the form

 $b_{l,s}(\theta) = \alpha_l \cos l\theta + \mathcal{O}(s), \qquad B_{l,s}(\theta) = \beta_l \cos l\theta + \mathcal{O}(s),$ 

with  $\alpha_l$  and  $\beta_l$  nonzero constants, such that the "Schiffer-type" problem can be solved in  $\Omega_{l,s}$ .

The result holds for example for l = 4 and for any l large enough.

# The shape of the domains $\Omega_{l,s}$



## The case l = 4



## Euler equations

#### **Euler equations**

The stationary incompressible Euler equations in the plane are:

$$\left\{ \begin{array}{ll} v\cdot\nabla v+\nabla p=0 & \mbox{ in } \mathbb{R}^2\,, \\ {\rm div}\,v=0 & \mbox{ in } \mathbb{R}^2\,. \end{array} \right.$$

Here  $v: \mathbb{R}^2 \to \mathbb{R}^2$  is the velocity of the fluid and  $p: \mathbb{R}^2 \to \mathbb{R}$  is the pressure.
#### Euler equations

The stationary incompressible Euler equations in the plane are:

$$\left\{ \begin{array}{ll} v\cdot\nabla v+\nabla p=0 & \mbox{ in } \mathbb{R}^2\,, \\ \mbox{ div } v=0 & \mbox{ in } \mathbb{R}^2\,. \end{array} \right.$$

Here  $v: \mathbb{R}^2 \to \mathbb{R}^2$  is the velocity of the fluid and  $p: \mathbb{R}^2 \to \mathbb{R}$  is the pressure.

Since div v = 0, there exists a "stream function" u with  $\nabla u = v^{\perp}$ .

#### **Euler equations**

The stationary incompressible Euler equations in the plane are:

$$\left\{ \begin{array}{ll} v\cdot\nabla v+\nabla p=0 & \mbox{ in } \mathbb{R}^2\,, \\ {\rm div}\,v=0 & \mbox{ in } \mathbb{R}^2\,. \end{array} \right.$$

Here  $v: \mathbb{R}^2 \to \mathbb{R}^2$  is the velocity of the fluid and  $p: \mathbb{R}^2 \to \mathbb{R}$  is the pressure.

Since div v = 0, there exists a "stream function" u with  $\nabla u = v^{\perp}$ .

**Old question.** Is it true that if (v, p) are  $C^1$  solutions and have compact support, then the streamlines (trajectories of the fluid) must be circular?

#### **Euler equations**

The stationary incompressible Euler equations in the plane are:

$$\left\{ \begin{array}{ll} v\cdot\nabla v+\nabla p=0 & \mbox{ in } \mathbb{R}^2\,, \\ {\rm div}\,v=0 & \mbox{ in } \mathbb{R}^2\,. \end{array} \right.$$

Here  $v: \mathbb{R}^2 \to \mathbb{R}^2$  is the velocity of the fluid and  $p: \mathbb{R}^2 \to \mathbb{R}$  is the pressure.

Since div v = 0, there exists a "stream function" u with  $\nabla u = v^{\perp}$ .

**Old question.** Is it true that if (v, p) are  $C^1$  solutions and have compact support, then the streamlines (trajectories of the fluid) must be circular?

This question is still open.

**Corollary.** There exist <u>continuous weak solutions</u> (v, p) with non-circular streamlines and compact support, and such that the stream function u satisfies an elliptic PDE.

**Corollary.** There exist <u>continuous weak solutions</u> (v, p) with non-circular streamlines and compact support, and such that the stream function u satisfies an elliptic PDE.

Take u as given by our main theorem. Then, define

$$v = \begin{cases} (\nabla u)^{\perp} & \text{in } \Omega, \\ 0 & \text{in } \mathbb{R}^2 \backslash \Omega, \end{cases} \quad p = \begin{cases} \mu \left( \frac{c_0^2}{2} - \frac{c_1^2}{2} \right) & \text{in } \Omega_{int}, \\ -\frac{1}{2}(|\nabla u|^2 + \mu(u^2 - c_0^2)) & \text{in } \Omega_{l,s}, \\ 0 & \text{in } \Omega_{ext}. \end{cases}$$

where  $\Omega_{int}$  and  $\Omega_{ext}$  are the two connected components of  $\mathbb{R}^2 \setminus \Omega_{l,s}$  respectively bounded and unbounded.

**Corollary.** There exist <u>continuous weak solutions</u> (v, p) with non-circular streamlines and compact support, and such that the stream function u satisfies an elliptic PDE.

Take u as given by our main theorem. Then, define

$$v = \begin{cases} (\nabla u)^{\perp} & \text{in } \Omega, \\ 0 & \text{in } \mathbb{R}^2 \backslash \Omega, \end{cases} \quad p = \begin{cases} \mu \left( \frac{c_0^2}{2} - \frac{c_1^2}{2} \right) & \text{in } \Omega_{int}, \\ -\frac{1}{2}(|\nabla u|^2 + \mu(u^2 - c_0^2)) & \text{in } \Omega_{l,s}, \\ 0 & \text{in } \Omega_{ext}. \end{cases}$$

where  $\Omega_{int}$  and  $\Omega_{ext}$  are the two connected components of  $\mathbb{R}^2 \setminus \Omega_{l,s}$  respectively bounded and unbounded.

**Gómez-Serrano, Park and Shi (Memoirs AMS, 2023).** Construction of other compactly supported  $C^0$  weak solutions (vortex-patch),

**Corollary.** There exist <u>continuous weak solutions</u> (v, p) with non-circular streamlines and compact support, and such that the stream function u satisfies an elliptic PDE.

Take u as given by our main theorem. Then, define

$$v = \begin{cases} (\nabla u)^{\perp} & \text{in } \Omega, \\ 0 & \text{in } \mathbb{R}^2 \backslash \Omega, \end{cases} \quad p = \begin{cases} \mu \left( \frac{c_0^2}{2} - \frac{c_1^2}{2} \right) & \text{in } \Omega_{int}, \\ -\frac{1}{2}(|\nabla u|^2 + \mu(u^2 - c_0^2)) & \text{in } \Omega_{l,s}, \\ 0 & \text{in } \Omega_{ext}. \end{cases}$$

where  $\Omega_{int}$  and  $\Omega_{ext}$  are the two connected components of  $\mathbb{R}^2 \setminus \Omega_{l,s}$  respectively bounded and unbounded.

**Gómez-Serrano, Park and Shi (Memoirs AMS, 2023).** Construction of other compactly supported  $C^0$  weak solutions (vortex-patch), but with stream function not satisfying any PDE.

**Crandall–Rabinowitz Theorem.** 

**Crandall–Rabinowitz Theorem.** X, Y = Banach spaces, U neighborhood of 0 in  $X, \Lambda$  real interval.

**Crandall–Rabinowitz Theorem.** X, Y = Banach spaces, U neighborhood of 0 in X,  $\Lambda$  real interval.  $F : \Lambda \times U \rightarrow Y C^1$  s.t.

**Crandall–Rabinowitz Theorem.** X, Y = Banach spaces, U neighborhood of 0 in X,  $\Lambda$  real interval.  $F : \Lambda \times U \rightarrow Y C^1$  s.t.

- i) F(a,0) = 0 for all  $a \in \Lambda$ ;
- ii) Ker  $D_w F(a_*, 0) = \mathbb{R} w_0$  for some  $a_* \in \Lambda$  and  $w_0 \in X \setminus \{0\}$ ;
- iii) codim Im  $D_w F(a_*, 0) = 1$ ;
- iv)  $D_a D_w F(a_*, 0)(w_0) \notin \operatorname{Im} D_w F(a_*, 0)$  ("transversality").

**Crandall–Rabinowitz Theorem.** X, Y = Banach spaces, U neighborhood of 0 in X,  $\Lambda$  real interval.  $F : \Lambda \times U \rightarrow Y C^1$  s.t.

- i) F(a,0) = 0 for all  $a \in \Lambda$ ;
- ii) Ker  $D_w F(a_*, 0) = \mathbb{R} w_0$  for some  $a_* \in \Lambda$  and  $w_0 \in X \setminus \{0\}$ ;
- iii) codim Im  $D_w F(a_*, 0) = 1$ ;
- iv)  $D_a D_w F(a_*, 0)(w_0) \notin \operatorname{Im} D_w F(a_*, 0)$  ("transversality").

If  $X = \mathbb{R} w_0 \oplus \dot{X}$ , there exists a nontrivial smooth curve

$$(-\epsilon,\epsilon) \to \dot{\Lambda} \times X, \quad s \mapsto (a(s), w(s))$$

such that

1) 
$$w(0) = 0$$
,  $a(0) = a_*$ , and  $s(w_0 + w(s)) \in U$ ;  
2)  $F(s(w_0 + w(s)), a(s)) = 0$ .

S. - 2010 (Calc. Var. PDEs). Bifurcation from the cylinder for the Serrin problem.

S. - 2010 (Calc. Var. PDEs). Bifurcation from the cylinder for the Serrin problem.

Del Pino, Pacard, Wei - 2015 (Duke). Bifurcation from domains bounded by unbounded CMC surfaces for the Serrin problem.

S. - 2010 (Calc. Var. PDEs). Bifurcation from the cylinder for the Serrin problem.

Del Pino, Pacard, Wei - 2015 (Duke). Bifurcation from domains bounded by unbounded CMC surfaces for the Serrin problem.

Ros, Ruiz, S. - 2020 (JEMS). Bifurcation from the exterior of a ball for the Serrin problem.

S. - 2010 (Calc. Var. PDEs). Bifurcation from the cylinder for the Serrin problem.

Del Pino, Pacard, Wei - 2015 (Duke). Bifurcation from domains bounded by unbounded CMC surfaces for the Serrin problem.

Ros, Ruiz, S. - 2020 (JEMS). Bifurcation from the exterior of a ball for the Serrin problem.

Kamburov, Sciaraffia - 2021 (Ann. IHP). Bifurcation from annuli for the Serrin problem (locally constant Dirichlet and Neumann).

S. - 2010 (Calc. Var. PDEs). Bifurcation from the cylinder for the Serrin problem.

Del Pino, Pacard, Wei - 2015 (Duke). Bifurcation from domains bounded by unbounded CMC surfaces for the Serrin problem.

Ros, Ruiz, S. - 2020 (JEMS). Bifurcation from the exterior of a ball for the Serrin problem.

Kamburov, Sciaraffia - 2021 (Ann. IHP). Bifurcation from annuli for the Serrin problem (locally constant Dirichlet and Neumann).

Fall, Minlend, Weth - 2023 (JEMS). Bifurcation from the cylinder for the Schiffer problem.

S. - 2010 (Calc. Var. PDEs). Bifurcation from the cylinder for the Serrin problem.

Del Pino, Pacard, Wei - 2015 (Duke). Bifurcation from domains bounded by unbounded CMC surfaces for the Serrin problem.

Ros, Ruiz, S. - 2020 (JEMS). Bifurcation from the exterior of a ball for the Serrin problem.

Kamburov, Sciaraffia - 2021 (Ann. IHP). Bifurcation from annuli for the Serrin problem (locally constant Dirichlet and Neumann).

Fall, Minlend, Weth - 2023 (JEMS). Bifurcation from the cylinder for the Schiffer problem.

Pacella, Ruiz, S. - 2024. Bifurcation from the the <u>relative</u> Serrin problem in a cylinder.

S. - 2010 (Calc. Var. PDEs). Bifurcation from the cylinder for the Serrin problem.

Del Pino, Pacard, Wei - 2015 (Duke). Bifurcation from domains bounded by unbounded CMC surfaces for the Serrin problem.

Ros, Ruiz, S. - 2020 (JEMS). Bifurcation from the exterior of a ball for the Serrin problem.

Kamburov, Sciaraffia - 2021 (Ann. IHP). Bifurcation from annuli for the Serrin problem (locally constant Dirichlet and Neumann).

Fall, Minlend, Weth - 2023 (JEMS). Bifurcation from the cylinder for the Schiffer problem.

Pacella, Ruiz, S. - 2024. Bifurcation from the the <u>relative</u> Serrin problem in a cylinder.

$$\Omega_a := \{ a < r < 1 \}.$$

 $\Omega_a := \{ a < r < 1 \}.$ 

 $\mu_2(a) =$  second radial Neumann eigenvalue on  $\Omega_a$ :

$$\Delta \psi_a + \mu_2(a)\psi_a = 0$$
 in  $\Omega_a$  ,  $|\nabla u| = 0$  on  $\partial \Omega_a$ 

 $\Omega_a := \{ a < r < 1 \}.$ 

 $\mu_2(a) =$  second radial Neumann eigenvalue on  $\Omega_a$ :

$$\Delta \psi_a + \mu_2(a)\psi_a = 0$$
 in  $\Omega_a$  ,  $|\nabla u| = 0$  on  $\partial \Omega_a$ 

We consider the canonical diffeomorphisme

$$\Phi_a^{b,B}: \Omega := \Omega_{\frac{1}{2}} \to \Omega_a^{b,B} := \{a + b(\theta) < r < 1 + B(\theta)\}$$

 $\Omega_a := \{ a < r < 1 \}.$ 

 $\mu_2(a) =$  second radial Neumann eigenvalue on  $\Omega_a$ :

$$\Delta \psi_a + \mu_2(a) \psi_a = 0$$
 in  $\Omega_a$  ,  $|\nabla u| = 0$  on  $\partial \Omega_a$ 

We consider the canonical diffeomorphisme

$$\Phi_a^{b,B} : \Omega := \Omega_{\frac{1}{2}} \to \Omega_a^{b,B} := \left\{ a + b(\theta) < r < 1 + B(\theta) \right\}.$$

and the pullback PDE:

$$L^{b,B}_a := (\Phi^{b,B}_a)^* \Delta + \mu_2(a) \, Id \ , \ \ L^{b,B}_a(u) = 0 \ \ {
m in} \ \ \Omega \, .$$

We need to impose symmetry.

We need to impose symmetry.

Let  $C_l^{k,\alpha}(\Omega)$  the space of  $C^{k,\alpha}(\Omega)$  functions with mode l, i.e. of type

 $\phi(r) \cos(l \theta)$ .

We need to impose symmetry.

Let  $C_l^{k,\alpha}(\Omega)$  the space of  $C^{k,\alpha}(\Omega)$  functions with mode l, i.e. of type  $\phi(r) \cos(l \theta)$ .

We need also regularity, i.e. very special Hölder spaces:

We need to impose symmetry.

Let  $C_l^{k,\alpha}(\Omega)$  the space of  $C^{k,\alpha}(\Omega)$  functions with mode l, i.e. of type  $\phi(r) \cos(l \theta)$ .

We need also regularity, i.e. very special Hölder spaces:

 $X^k := \{ u \in C^{k,\alpha}_l(\Omega) \, | \, \partial_r u \in C^{k,\alpha}(\Omega) \}.$ 

We need to impose symmetry.

Let  $C_l^{k,\alpha}(\Omega)$  the space of  $C^{k,\alpha}(\Omega)$  functions with mode l, i.e. of type  $\phi(r) \cos(l \theta)$ .

We need also regularity, i.e. very special Hölder spaces:

$$X^k := \{ u \in C^{k,\alpha}_l(\Omega) \, | \, \partial_r u \in C^{k,\alpha}(\Omega) \}.$$

 $X_D^k:=\{u\in X^k\,|\, u=0\, {\rm on}\, \partial\Omega\}.$ 

We need to impose symmetry.

Let  $C_l^{k,\alpha}(\Omega)$  the space of  $C^{k,\alpha}(\Omega)$  functions with mode l, i.e. of type  $\phi(r) \cos(l \theta)$ .

We need also regularity, i.e. very special Hölder spaces:

$$X^k := \{ u \in C^{k,\alpha}_l(\Omega) \, | \, \partial_r u \in C^{k,\alpha}(\Omega) \}.$$

 $X_D^k := \{ u \in X^k \, | \, u = 0 \text{ on } \partial \Omega \}.$ 

$$X_{DN}^k := \{ u \in X^k \mid u = 0, \partial_r u = 0 \text{ on } \partial \Omega \}.$$

We need to impose symmetry.

Let  $C_l^{k,\alpha}(\Omega)$  the space of  $C^{k,\alpha}(\Omega)$  functions with mode l, i.e. of type  $\phi(r) \cos(l \theta)$ .

We need also regularity, i.e. very special Hölder spaces:

$$\begin{split} X^{k} &:= \{ u \in C_{l}^{k,\alpha}(\Omega) \mid \partial_{r} u \in C^{k,\alpha}(\Omega) \}. \\ X^{k}_{D} &:= \{ u \in X^{k} \mid u = 0 \text{ on } \partial\Omega \}. \\ X^{k}_{DN} &:= \{ u \in X^{k} \mid u = 0, \partial_{r} u = 0 \text{ on } \partial\Omega \}. \\ Y &:= C_{l}^{1,\alpha}(\Omega) + X^{0}_{D} , \quad \|y\|_{Y} = \inf\{\|y_{1}\| + \|y_{2}\| \mid, y = y_{1} + y_{2} \}. \end{split}$$

The operator to performe the bifurcation

## The operator to performe the bifurcation

For  $v \in X_D^2$  small we define two functions on  $\partial \Omega$ 

$$b_v(\theta) = c_1(a)\partial_r v\left(\frac{1}{2}, \theta\right) , \quad B_v(\theta) = c_2(a)\partial_r v(1, \theta) ,$$

for some suitable constants  $c_1(a), c_2(a)$ .
## The operator to performe the bifurcation

For  $v \in X_D^2$  small we define two functions on  $\partial \Omega$ 

$$b_v(\theta) = c_1(a)\partial_r v\left(\frac{1}{2}, \theta\right) , \quad B_v(\theta) = c_2(a)\partial_r v(1, \theta) ,$$

for some suitable constants  $c_1(a), c_2(a)$ . We define also

$$w_{v}(r,\theta) = v(r,\theta) + \frac{\overline{\psi}_{a}'(r)}{2(1-a)} \Big[ 2(1-r)b_{v}(\theta) + (2r-1)B_{v}(\theta) \Big] \in \mathcal{X}_{DN}^{2,\alpha}$$

## The operator to performe the bifurcation

For  $v \in X_D^2$  small we define two functions on  $\partial \Omega$ 

$$b_v(\theta) = c_1(a)\partial_r v\left(\frac{1}{2}, \theta\right) , \quad B_v(\theta) = c_2(a)\partial_r v(1, \theta) ,$$

for some suitable constants  $c_1(a), c_2(a)$ . We define also

$$w_{v}(r,\theta) = v(r,\theta) + \frac{\overline{\psi}_{a}'(r)}{2(1-a)} \Big[ 2(1-r)b_{v}(\theta) + (2r-1)B_{v}(\theta) \Big] \in \mathcal{X}_{DN}^{2,\alpha}.$$

With this, we define the operator:

$$F(a,v) := L_a^{b_v, B_v}[\bar{\psi}_a + w_v] \in Y$$

where  $\bar{\psi}_a$  is the pullback of  $\psi_a$  for the operator  $L_a^{0,0}$ .

# The linearized operator

The linearized operator

Proposition. We have

$$D_w F(a, 0) := L_a^{0,0} : X_D^2 \to Y$$

The linearized operator

Proposition. We have

$$D_w F(a, 0) := L_a^{0,0} : X_D^2 \to Y$$

(in particular it is a Fredholm operator of index 0).

Proposition. We have

$$D_w F(a, 0) := L_a^{0,0} : X_D^2 \to Y$$

(in particular it is a Fredholm operator of index 0). And it becomes degenerate when  $\mu_2(a)$  is also a Dirichlet eigenvalue.

Proposition. We have

$$D_w F(a,0) := L_a^{0,0} : X_D^2 \to Y$$

(in particular it is a Fredholm operator of index 0). And it becomes degenerate when  $\mu_2(a)$  is also a Dirichlet eigenvalue.

We denote by  $\lambda_l(a)$  the first Dirichlet eigenvalue of  $\Omega_a$  of mode l (i.e. with eigenfunction of the form  $\phi(r) \cos(l\theta)$ ).

Proposition. We have

$$D_w F(a,0) := L_a^{0,0} : X_D^2 \to Y$$

(in particular it is a Fredholm operator of index 0). And it becomes degenerate when  $\mu_2(a)$  is also a Dirichlet eigenvalue.

We denote by  $\lambda_l(a)$  the first Dirichlet eigenvalue of  $\Omega_a$  of mode l (i.e. with eigenfunction of the form  $\phi(r) \cos(l\theta)$ ).

#### Proposition. We have:

- If *a* is close to 1, then  $\mu_2(a) > \lambda_l(a)$
- If a is close to 0 and  $l \ge 4$ , then  $\mu_2(a) < \lambda_l(a)$

So there exist  $a = a_*$  such that the two eigenvalues are equal, and there  $D_w F(a, 0)$  is degenerate.

In order to apply the Crandall-Rabinowitz Bifurcatiion Theorem we just need the "transversality condition".

In order to apply the Crandall-Rabinowitz Bifurcatiion Theorem we just need the "transversality condition".

Such condition in our case is

 $\mu_2'(a_*) \neq \lambda_l'(a_*)$ 

In order to apply the Crandall-Rabinowitz Bifurcatiion Theorem we just need the "transversality condition".

Such condition in our case is

 $\mu_2'(a_*) \neq \lambda_l'(a_*)$ 

**Proposition.** For l = 4 and for any integer l large enough we have

 $\mu_2'(a_*) > \lambda_l'(a_*)$ 

In order to apply the Crandall-Rabinowitz Bifurcatiion Theorem we just need the "transversality condition".

Such condition in our case is

 $\mu_2'(a_*) \neq \lambda_l'(a_*)$ 

**Proposition.** For l = 4 and for any integer l large enough we have

 $\mu_2'(a_*) > \lambda_l'(a_*)$ 

 $\Longrightarrow$  all the hypothesis of the Crandall-Rabinowitz theorem

Mostly maximum principle

# THANK YOU FOR YOUR ATTENTION